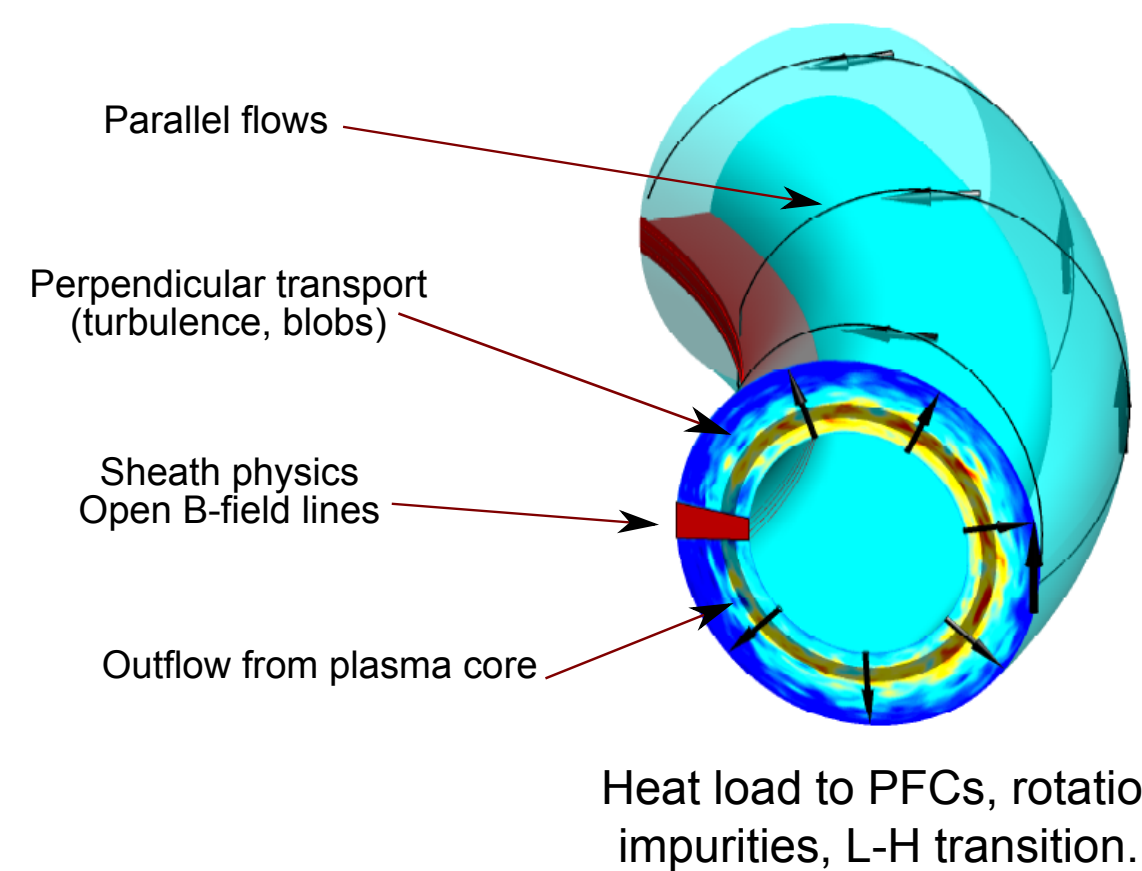


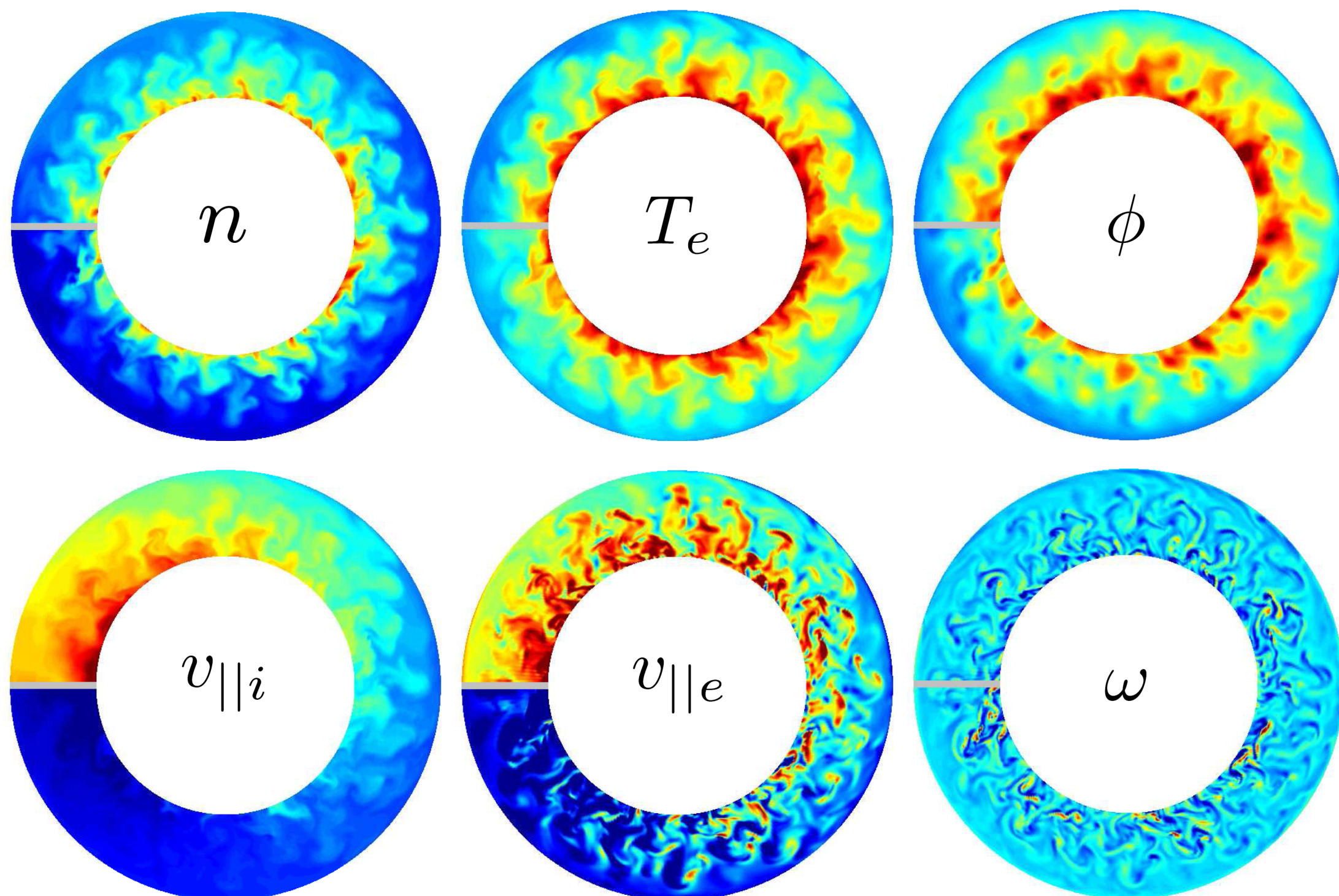
Introduction



- In the tokamak scrape-off layer (SOL) magnetic field lines are open, **channeling heat onto device wall**
- **Large computational effort** devoted to understand width of the heat flux channel using **Global Braginskii Solver (GBS)** plasma turbulence code
- SOL dynamics studied via **direct numerical simulations** capable of resolving **turbulent dynamics at experimental parameters**
- **Latest numerical developments in GBS**
 - Fast, efficient discretization of parallel dynamics
 - Matrix-free parallel multigrid solver for the Poisson equation
- Using new numerical approach, simulations of even larger tokamaks are possible using $\sim 10^4$ CPUs

GBS, a tool to study SOL turbulent plasma dynamics

- Flux-driven plasma turbulence code to study SOL heat and particle transport
- Code fully verified using **method of manufactured solutions** [Riva *et al.*, PoP (2014)]
- Spatial discretization using finite differences, RK4 integration in time
- Parallelized using domain decomposition in x and z axes
- For medium tokamak size $\rho_*^{-1} \approx 2000$, **excellent parallel scalability up to 1024 cores**



Drift-reduced fluid equations for plasma turbulence

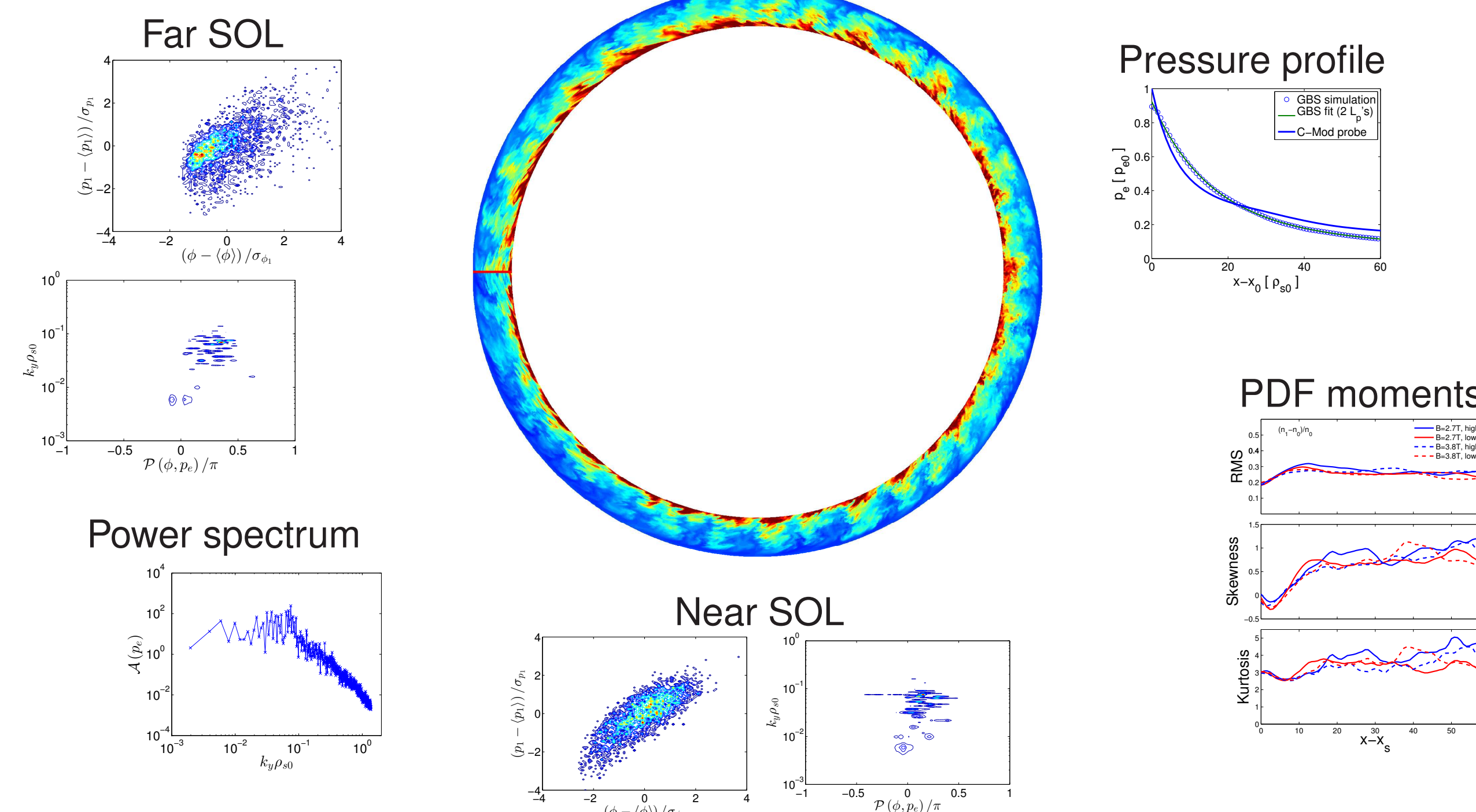
- Low-frequency, collisional, electromagnetic turbulence driven by plasma gradients
- Large fluctuations $\mathcal{O}(1)$, no length scale separation
- Drift-reduced Braginskii eqns with orderings $k_{\perp} \gg k_{\parallel}$, $d/dt \ll \omega_{ci}$ [Ricci *et al.*, PPCF 2012]:

$$\begin{aligned} \frac{\partial n}{\partial t} &= -\rho_*^{-1}[\phi, n] + \frac{2}{B}[C(p_e) - nC(\phi)] - \nabla_{\parallel} \cdot (n\mathbf{v}_{||e}) + S_n \\ \frac{\partial \nabla_{\perp}^2 \phi}{\partial t} &= -\rho_*^{-1}[\phi, \nabla_{\perp}^2 \phi] - \mathbf{v}_{||i} \nabla_{\parallel} \nabla_{\perp}^2 \phi + \frac{B^2}{n} \nabla_{\parallel} j_{||} + \frac{2B}{n} C(p) \\ \frac{\partial \chi}{\partial t} &= -\rho_*^{-1}[\phi, \mathbf{v}_{||e}] - \mathbf{v}_{||e} \nabla_{\parallel} \mathbf{v}_{||e} + \frac{m_i}{m_e} \left(\nu \frac{j_{||}}{n} + \nabla_{\parallel} \phi - \frac{1}{n} \nabla_{\parallel} p_e - 0.71 \nabla_{\parallel} T_e \right) \\ \frac{\partial \mathbf{v}_{||i}}{\partial t} &= -\rho_*^{-1}[\phi, \mathbf{v}_{||i}] - \mathbf{v}_{||i} \nabla_{\parallel} \mathbf{v}_{||i} - \frac{1}{n} \nabla_{\parallel} p \\ \frac{\partial T_e}{\partial t} &= -\rho_*^{-1}[\phi, T_e] - \mathbf{v}_{||e} \nabla_{\parallel} T_e + \frac{4}{3} \frac{T_e}{B} \left[\frac{1}{n} C(p_e) + \frac{5}{2} C(T_e) - C(\phi) \right] + \frac{2}{3} T_e [0.71 \nabla_{\parallel} j_{||} - \nabla_{\parallel} \mathbf{v}_{||e}] + S_{T_e} \\ \frac{\partial T_i}{\partial t} &= -\rho_*^{-1}[\phi, T_i] - \mathbf{v}_{||i} \nabla_{\parallel} T_i + \frac{4}{3} \frac{T_i}{B} \left[C(T_e) + \frac{T_e}{n} C(n) - C(\phi) \right] + \frac{2}{3} T_i (\mathbf{v}_{||i} - \mathbf{v}_{||e}) \frac{\nabla_{\parallel} n}{n} - \frac{2}{3} T_i \nabla_{\parallel} \mathbf{v}_{||e} - \frac{10}{3} \frac{T_i}{B} C(T_i) \\ \nabla_{\perp}^2 \psi &= n(\mathbf{v}_{||i} - \mathbf{v}_{||e}) = j_{||}, \quad \chi = \mathbf{v}_{||e} + \frac{m_i \beta}{m_e 2} \psi, \quad \rho_* = \rho_e / R \\ \nabla_{\parallel} f &= \mathbf{b}_0 \cdot \nabla f + \rho_*^{-1} \frac{\beta}{2} [\psi, f] \end{aligned}$$

- Normalized units used throughout: $L_{\perp} \rightarrow \rho_s$, $L_{\parallel} \rightarrow R$, $t \rightarrow R/c_s$, $\nu = ne^2 c_s / (m_i \sigma_{\parallel} R)$

SOL turbulent dynamics revealed through computer simulations

- GBS simulations at **realistic parameters** reproduce features found in SOL of tokamak discharges
 - Pressure profile is a decaying exponential with **2 decay lengths**
 - Large fluctuations $\sim 30\%$, skewed PDF shows **presence of blobs**
 - **Two distinct regions**: drift dominated vs. interchange dominated
- **Detailed comparison** between code and experimental measurements in progress
 - Dedicated experiments carried out at Alcator C-Mod (MIT) using **state of the art diagnostics**

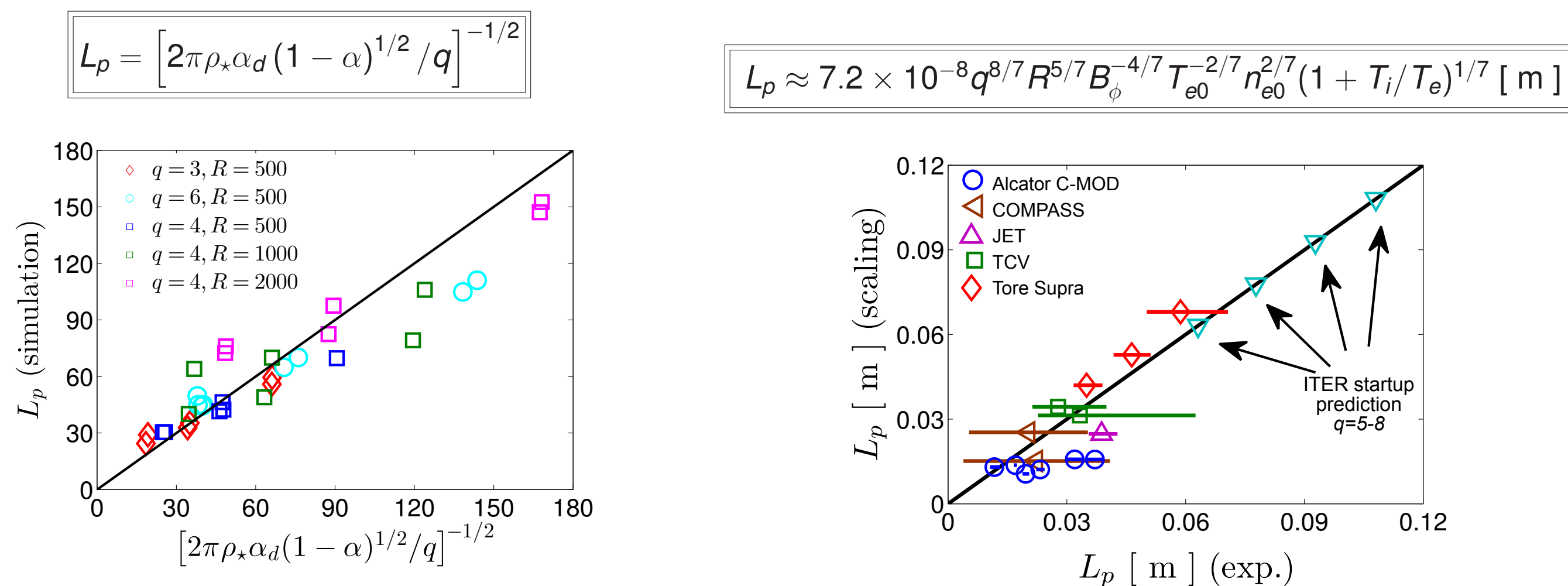


- Analysis of a large simulation scan revealed turbulent saturation mechanism, non-linear instability regimes, equilibrium electric field, effects of parallel dynamics...
- SOL width can be calculated **analytically** by considering *gradient removal saturated turbulence*:

$$L_p = q \frac{\gamma_b}{k_b} \quad \gamma_b = \sqrt{2/(\rho_* L_p)} \quad k_b = \sqrt{(1-\alpha)/(\nu \gamma_b)}/q$$

$$\alpha_d^{-1} = 2^{7/4} \nu^{1/2} (\rho_* L_p)^{1/4} \pi q \quad \alpha = \frac{q^2 \beta}{\rho_* L_p}$$

- Dimensionless and engineering parameter scalings of the SOL width follow [Halpern *et al.*, NF 2013]:



- Scalings obtained from least-squares-fitting of all simulation data verify our theory:

$$L_p = 0.42 q^{0.55} \rho_*^{-0.53} \alpha_d^{-0.32} (1-\alpha)^{-0.24} \rightarrow L_p \sim q^{0.98} R^{0.63} B^{-0.56} [\text{m}]$$

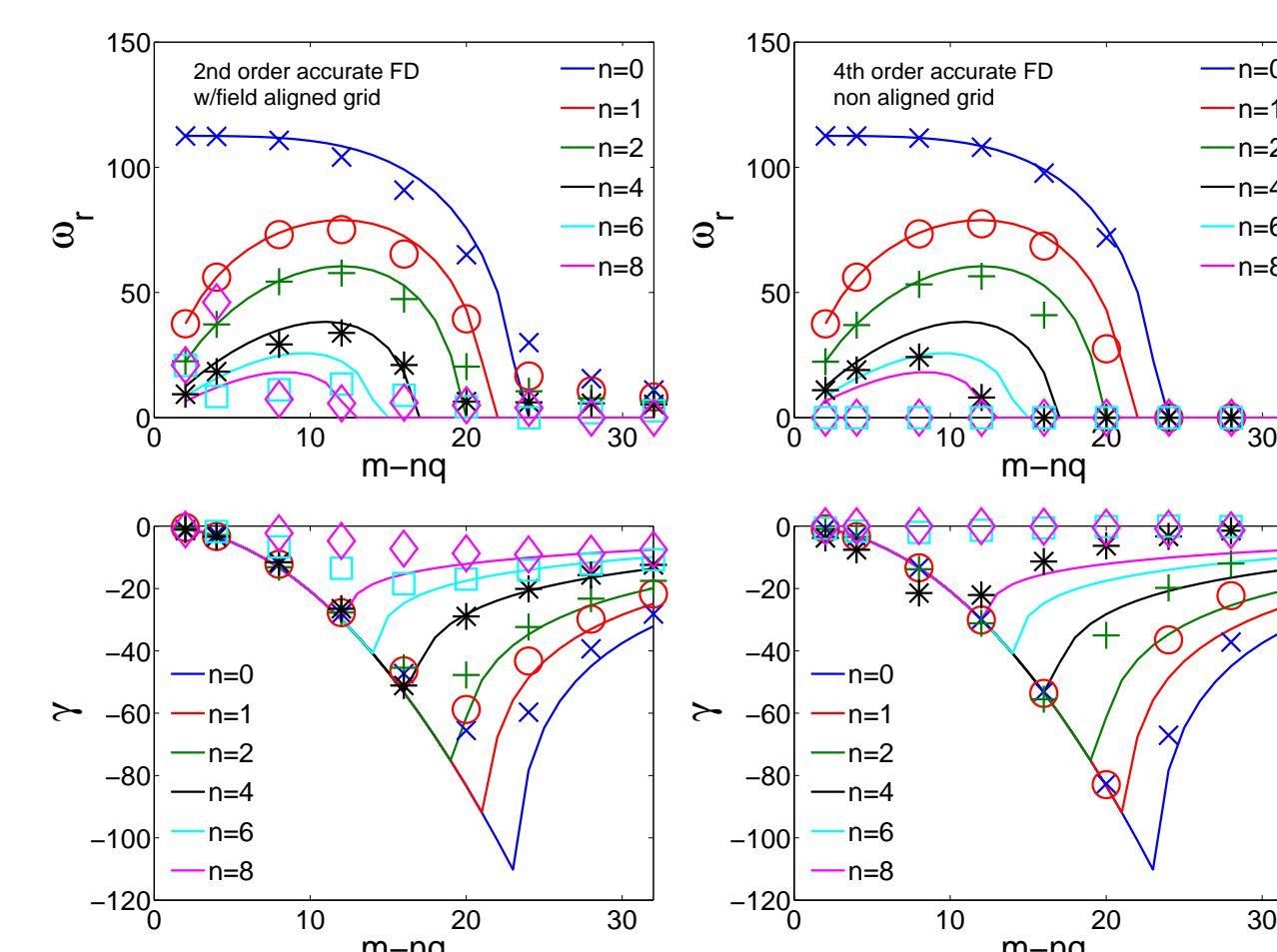
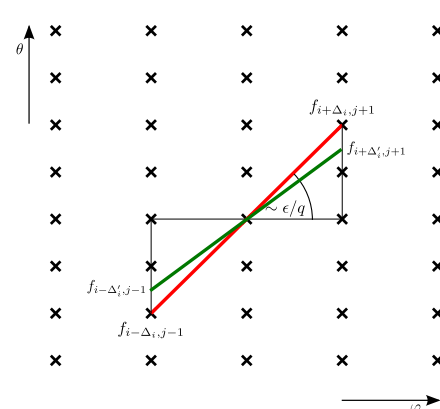
Recent numerical developments

Parallel dynamics discretization schemes

- Description of parallel dynamics **essential** for efficient and stable plasma turbulence codes
- Turbulence strongly anisotropic, elongated turbulent structures aligned to B -field
- Basic wave phenomena regulating plasma dynamics \rightarrow shear-Alfvén waves
- **We test field-aligned and non-field aligned schemes to find optimal algorithm for GBS**
- Study parallel dynamics using simplified model
 - Evaluate numerical ω , γ using simple code:

$$\frac{\partial \nabla_{\perp}^2 \phi}{\partial t} = -\nabla_{\parallel} \mathbf{v}_{||e} \cdot \nabla_{\perp} \phi, \quad \frac{\partial \mathbf{v}_{||e}}{\partial t} = \frac{m_i}{m_e} \left(\nu \frac{j_{||}}{n} + \nabla_{\parallel} \phi - \frac{1}{n} \nabla_{\parallel} p_e - 0.71 \nabla_{\parallel} T_e \right)$$

- Simple analytical solution $\omega^2 = \omega_0^2 - \gamma^2$, $\omega_0 = \sqrt{m_i/m_e} k_{\parallel}/k_{\perp}$, $\gamma = 2\eta k_{\parallel}^2/3$
- Evolve perturbations $\sim \sin(my - nz)$ using different numerical schemes



Field aligned and 4th order FD schemes reproduce parallel dynamics accurately and efficiently

Parallel, matrix-free multigrid solver

Efficient, **scalable** inversion of **time dependent** operators, e.g.:

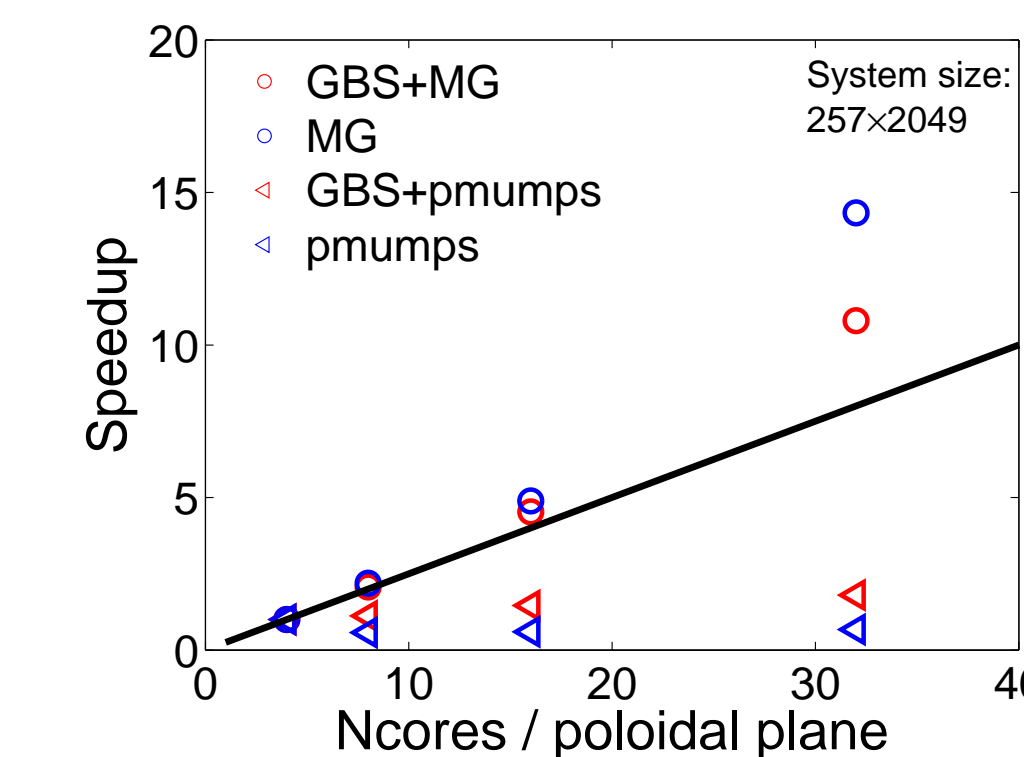
- Electromagnetic Ohm's law

$$\left[\nabla_{\perp}^2 - \frac{\beta m_i}{2 m_e} n \right] \mathbf{v}_{||e} = S_{v_{||e}}$$

- Non-Boussinesq polarization equation

$$[\nabla_{\perp}^2 + \nabla_{\perp} \ln n \cdot \nabla_{\perp}] \phi = S_{\phi}$$

- **Avoid LU decomposition in every time step**

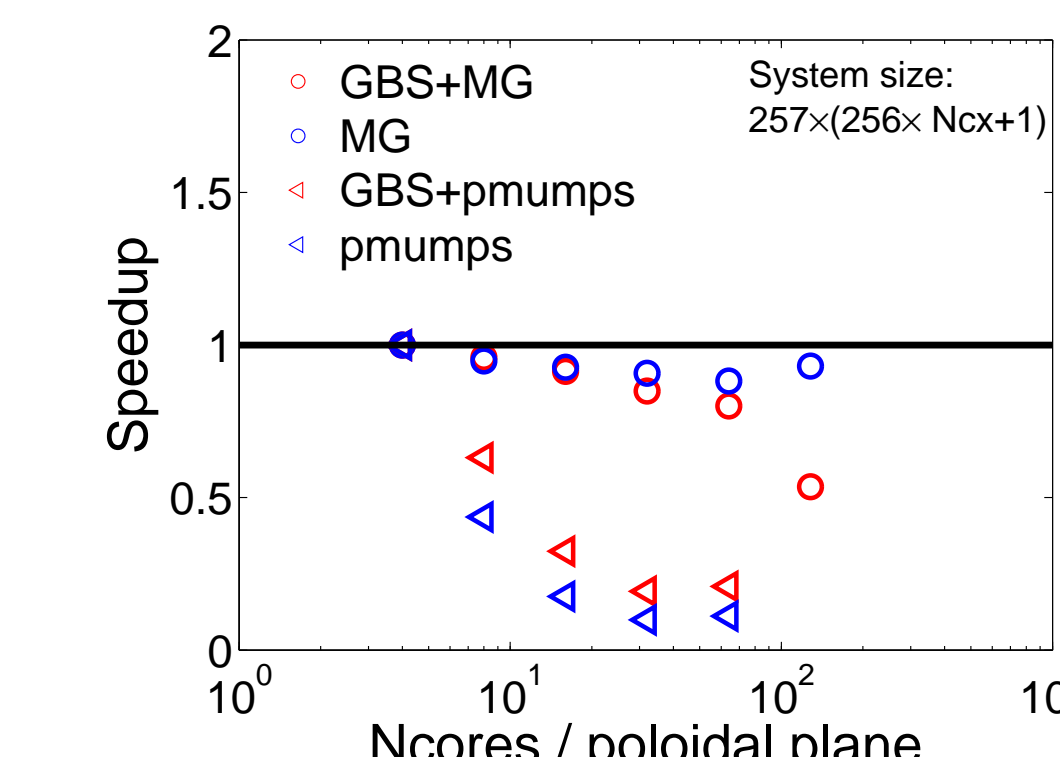


Developed multigrid solver in 2D Cartesian grid mapped to 2D domain decomposition

- FD matrix partitioned using **stencil notation**

- Full weighting restriction, **bilinear interpolation** for prolongation, damped Jacobi relaxation, $\Omega = 0.9$

- Carry out simple test with $\nabla_{\perp}^2 \phi = \omega$, **weak and strong scalings** in MonteRosa shown below



Conclusions

- Numerical studies of tokamak turbulence improved understanding of plasma-wall interaction
- Simulations show many features found in experimental measurements
- Developed predictive theory for SOL width of inner-wall limited tokamak plasmas
- Analyzed wave dispersion and damping with numerical schemes appropriate for X-point geometry
- Demonstrated super-linear GBS speedup using parallel multigrid solver

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Outlook and GBS development plans

- Validate SOL model against dedicated experiments in Alcator C-Mod
- Parallelize GBS in y axis \rightarrow factor 10 speedup
- Use several numerical approaches to model magnetic separatrix / X-point
- Port GBS to OpenMP/MPI hybrid, then to GPU architecture
- Carry out kinetic simulations of plasma biasing \rightarrow turbulence control